



# **Relativistic Electrodynamics**

**M.Sc. 2<sup>nd</sup> Semester**

**MPHYCC-6: Electrodynamics and Plasma Physics**

**Unit IV (Part 1)**

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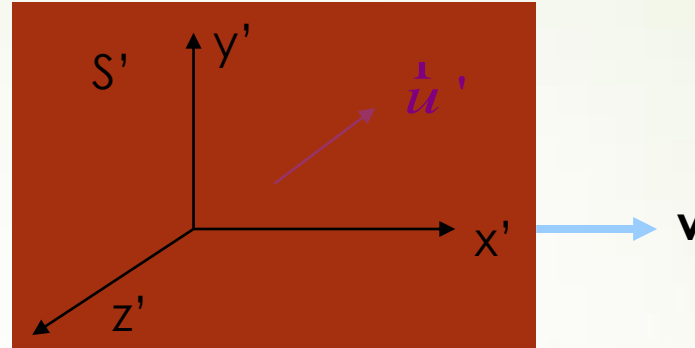
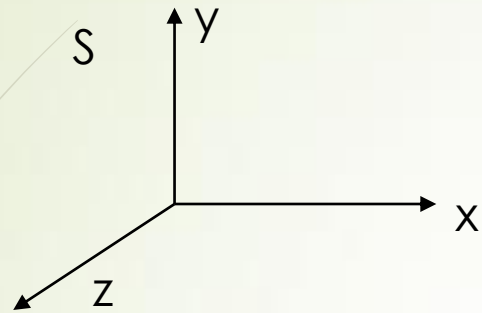
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# Lecture Outline

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- Special Theory of Relativity
- Transformation of electric force
- 4-vector notations
- Covariant 4-vector and contravariant 4-vector
- Covariant form of the electromagnetic field
- Covariant form of the Maxwell equations
  
- Assignment

# Special theory of relativity



space-time

$$x = (x_0, \mathbf{x}) = (ct, \mathbf{x})$$

$$x_0 = \gamma(x'_0 + \beta x'_1)$$

$$x_1 = \gamma(x'_1 + \beta x'_0)$$

$$x_2 = x'_2$$

$$x_3 = x'_3$$

$$\beta \equiv v/c, \quad \gamma \equiv \frac{1}{\sqrt{1 - (v/c)^2}}$$

energy-momentum

$$p = (p_0, \mathbf{p}) = (E/c, \mathbf{p})$$

$$E = \gamma_u mc^2, \quad \mathbf{p} = \gamma_u m \mathbf{u}$$

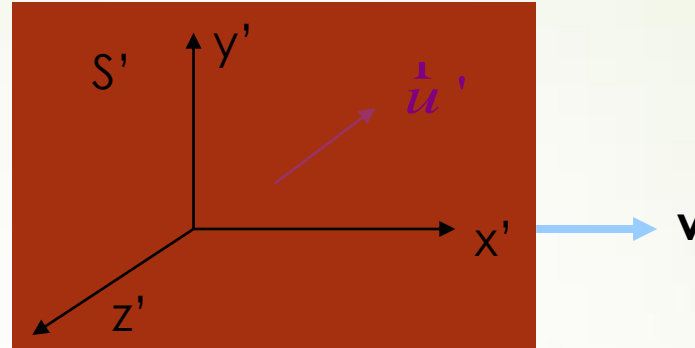
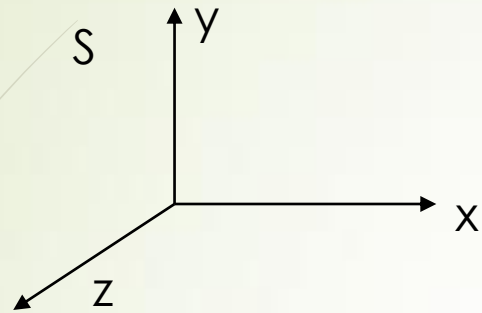
$$p_0 = \gamma_v (p'_0 + \beta p'_1)$$

$$p_1 = \gamma_v (p'_1 + \beta p'_0)$$

$$p_2 = p'_2$$

$$p_3 = p'_3$$

# STR



velocity,  $\dot{\mathbf{x}} = d\mathbf{x} / dt$

$$u_1 = \frac{u'_1 + v}{1 + \frac{u'_1 v}{c^2}}$$

$$u_2 = \frac{u'_2}{\gamma \left( 1 + \frac{u'_1 v}{c^2} \right)}$$

$$u_3 = \frac{u'_3}{\gamma \left( 1 + \frac{u'_1 v}{c^2} \right)}$$

force,  $\dot{\mathbf{r}} = d\mathbf{p} / dt$

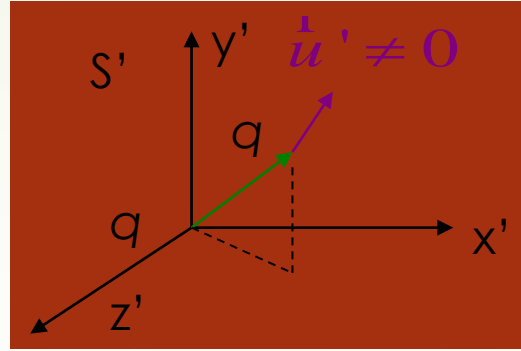
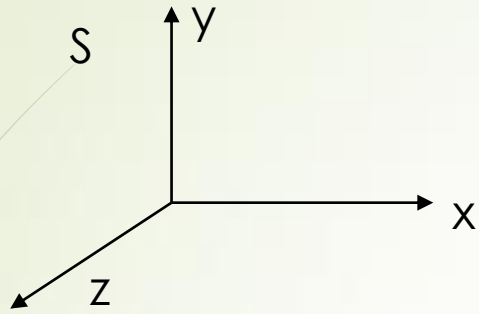
$$F_1 = F'_1 + \frac{v}{c^2 + u'_1 v} (u'_2 F'_2 + u'_3 F'_3)$$

$$F_2 = \frac{F'_2}{\gamma \left( 1 + \frac{u'_1 v}{c^2} \right)}$$

$$F_3 = \frac{F'_3}{\gamma \left( 1 + \frac{u'_1 v}{c^2} \right)}$$

$$\left( \frac{dE'}{dt'} = \mathbf{F}' \cdot \mathbf{u}' \text{ has been used} \right)$$

# Transformation of electric force



$$\hat{r}' = \gamma(x - vt)\hat{x} + y\hat{y} + z\hat{z}$$

$$\mathbf{v}$$

$$\hat{r} \equiv (x - vt)\hat{x} + y\hat{y} + z\hat{z}$$

$$F'_x = \frac{q^2 x'}{r'^3}, F'_y = \frac{q^2 y'}{r'^3}, F'_z = \frac{q^2 z'}{r'^3}$$

$$\Rightarrow F_x = F'_x + \frac{v}{c^2 + u'_x v} (u'_y F'_y + u'_z F'_z)$$

$$\text{K} = \frac{e^2 \gamma}{r'^3} \left( x - vt + \frac{u_y v}{c^2} y + \frac{u_z v}{c^2} z \right)$$

$$F_y = \frac{F'_y}{\gamma \left( 1 + \frac{u'_x v}{c^2} \right)} = \frac{e^2 \gamma y}{r'^3} \left( 1 - \frac{u_x v}{c^2} \right)$$

$$F_z = \frac{F'_z}{\gamma \left( 1 + \frac{u'_x v}{c^2} \right)} = \frac{e^2 \gamma z}{r'^3} \left( 1 - \frac{u_x v}{c^2} \right)$$

Vector notation

$$\hat{r} = \frac{q^2 \gamma}{r'^3} \left\{ \hat{r} + \frac{v}{c^2} \left[ (u_y y + u_z z) \hat{x} - u_x (y\hat{y} + z\hat{z}) \right] \right\}$$

$$= \frac{q^2 \gamma}{r'^3} \left\{ \hat{r} + \frac{1}{c^2} \hat{u} \times (\hat{v} \times \hat{r}) \right\}$$

$$= q \left( \hat{E} + \frac{\hat{u}}{c} \times \hat{B} \right)$$

where  $\hat{E} = \frac{q\gamma}{r'^3} \hat{r}$

$$\hat{B} \equiv \frac{q\gamma}{r'^3} \left( \frac{\hat{v}}{c} \times \hat{r} \right) = \frac{\hat{v}}{c} \times \hat{E}$$

( $\hat{u} = \hat{v}$  on previous page.)

# 4-vector notations

- Space-time ( $ct, \mathbf{x}$ )
- energy-momentum ( $E/c, \mathbf{p}$ )
- scalar potential-vector potential ( $\phi, \mathbf{A}$ )
- 4-dim space-time operator (below)
- charge density-charge current (below)
- ...

but not the usual velocity, acceleration, force...

4-dim space-time operator

$$\frac{\partial}{\partial x} = \left( \frac{\partial}{\partial x_0}, \frac{\partial}{\partial \mathbf{x}} \right) \text{ or } (\partial_0, \partial_{\mathbf{x}})$$

$$\partial_0 = \gamma (\partial'_0 - \beta \partial'_1)$$

$$\partial_1 = \gamma (\partial'_1 - \beta \partial'_0)$$

$$\partial_2 = \partial'_2$$

$$\partial_3 = \partial'_3$$

charge density-charge current

$$j = (j_0, \mathbf{j}) = (c\rho, \mathbf{j})$$

$$dq = \rho dV, \mathbf{j} = \rho \mathbf{u}$$

$$j_0 = \gamma_v (j'_0 + \beta j'_1)$$

$$j_1 = \gamma_v (j'_1 + \beta j'_0)$$

$$j_2 = j'_2$$

$$j_3 = j'_3$$

# Covariant 4-vector and contravariant 4-vector

Vectors that transform like  $x$  are called **contravariant** vectors; Vectors that transform like  $\partial / \partial x$  are called **covariant** vectors. Their notations are distinguished by the position of the index.

$$dx^{\alpha \uparrow} = \frac{\partial x^{\alpha}}{\partial x'^{\beta}} dx'^{\beta}$$

$$\frac{\partial}{\partial x^{\alpha}} = \frac{\partial x'^{\beta}}{\partial x^{\alpha}} \frac{\partial}{\partial x'^{\beta}} \quad \text{or} \quad \partial_{\alpha \downarrow} = \frac{\partial x'^{\beta}}{\partial x^{\alpha}} \partial'_{\beta}$$

$$\begin{pmatrix} \gamma & \gamma\beta & 0 & 0 \\ \gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \xleftrightarrow{\text{inverse}} \begin{pmatrix} \gamma & -\gamma\beta & 0 & 0 \\ -\gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

Raising and lowering of the index

$$A_{\alpha} = g_{\alpha\beta} A^{\beta}$$

$$A^{\alpha} = g^{\alpha\beta} A_{\beta}$$

$$\text{where } g_{\alpha\beta} = g^{\alpha\beta} \equiv \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

$g$  (called a metric tensor) converts a contravariant vector to a covariant vector, and vice versa

# Inner product between two 4-vectors

$$A^\alpha = (A^0, A^i), B^\alpha = (B^0, B^i)$$

$$A \cdot B \equiv A^\alpha B_\alpha = A_\alpha B^\alpha = A^\alpha g_{\alpha\beta} B^\beta = A_\alpha g^{\alpha\beta} B_\beta$$

$$= A^0 B_0 + A^i B_i$$

$$= A^0 B^0 - A^i B^i$$

The inner product is invariant under Lorentz transformation because  $A^\alpha$  and  $B_\alpha$  transform oppositely.

## Relativistic invariants

$$x^2 = c^2 t^2 - |\mathbf{x}|^2$$

$$p^2 = (E/c)^2 - |\mathbf{p}|^2 = m^2 c^2$$

$$\partial_\alpha j^\alpha = \frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{j} = 0$$

... etc

Conversely, any linear transformation that leaves  $x^2$  (or other inner product) invariant must be a Lorentz transformation (including spatial rotation).

Analogy: any linear transformation that leaves  $|\mathbf{x}|^2$  (or other inner product) invariant must be a rotation.



# Covariant form of the electromagnetic field

$$\vec{E} = -\frac{1}{c} \frac{\partial \vec{A}}{\partial t} - \nabla \phi; \quad \vec{B} = \nabla \times \vec{A}$$

can be written as

$$E_i = -(\partial^0 A^i - \partial^i A^0); \quad \varepsilon^{ijk} B_k = -(\partial^i A^j - \partial^j A^i)$$

it's natural to define

$$F^{\alpha\beta} = \partial^\alpha A^\beta - \partial^\beta A^\alpha$$

$$= \begin{pmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & -B_z & B_y \\ E_y & B_z & 0 & -B_x \\ E_z & -B_y & B_x & 0 \end{pmatrix}$$

$$F_{\alpha\beta} \equiv \partial_\alpha A_\beta - \partial_\beta A_\alpha$$

$$= \begin{pmatrix} 0 & E_x & E_y & E_z \\ -E_x & 0 & -B_z & B_y \\ -E_y & B_z & 0 & -B_x \\ -E_z & -B_y & B_x & 0 \end{pmatrix}$$

$$B_k = -\frac{1}{2} \varepsilon_{ijk} F^{ij}$$

Transformation of the field strength tensor

$$F^{\alpha\beta} = \frac{\partial x^\alpha}{\partial x'^\nu} \frac{\partial x'^\beta}{\partial x'^\delta} F'^{\nu\delta}$$

$$\begin{pmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & -B_z & B_y \\ E_y & B_z & 0 & -B_x \\ E_z & -B_y & B_x & 0 \end{pmatrix} = \begin{pmatrix} \gamma & \gamma\beta & 0 & 0 \\ \gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & -E'_x & -E'_y & -E'_z \\ E'_x & 0 & -B'_z & B'_y \\ E'_y & B'_z & 0 & -B'_x \\ E'_z & -B'_y & B'_x & 0 \end{pmatrix} \begin{pmatrix} \gamma & \gamma\beta & 0 & 0 \\ \gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

or

# Covariant form of the Maxwell equations

$$\nabla \cdot \overset{\mathbf{r}}{E} = 4\pi\rho$$

$$\nabla \cdot \overset{\mathbf{r}}{B} = 0$$

$$\nabla \times \overset{\mathbf{r}}{E} + \frac{1}{c} \frac{\partial \overset{\mathbf{r}}{B}}{\partial t} = 0$$

$$\nabla \times \overset{\mathbf{r}}{B} - \frac{1}{c} \frac{\partial \overset{\mathbf{r}}{E}}{\partial t} = \frac{4\pi}{c} \overset{\mathbf{r}}{J}$$



$$\partial_\alpha F^{\alpha\beta} = \frac{4\pi}{c} J^\beta \quad \Rightarrow \quad \partial_\alpha J^\alpha = 0$$

$$\partial^\alpha F^{\beta\gamma} + \partial^\beta F^{\gamma\alpha} + \partial^\gamma F^{\alpha\beta} = 0$$

(Bianchi identity)

## Covariant form of the Lorentz force equation

The usual velocity  $\mathbf{v}$  is not part of a 4-vector since  $t$  is not invariant under the Lorentz transformation

4-velocity  $u^\alpha \equiv \frac{dx^\alpha}{d\tau}$

where  $d\tau$  is a relativistic inv. w/ the dimension of time

choose  $d\tau = \frac{ds}{c}$  ( $ds^2 = c^2 dt^2 - |d\overset{\mathbf{r}}{x}|^2$ )

$$= dt \sqrt{1 - \beta^2} = \frac{dt}{\gamma} \quad (=dt' \text{ when } \overset{\mathbf{r}}{x}' = 0)$$

$$u = (\gamma c, \gamma \overset{\mathbf{r}}{v})$$

$$\frac{d\overset{\mathbf{r}}{p}}{dt} = q \left( \overset{\mathbf{r}}{E} + \frac{\overset{\mathbf{r}}{v}}{c} \times \overset{\mathbf{r}}{B} \right)$$

$$\Leftrightarrow \frac{dp^i}{d\tau} = \frac{q}{c} F^{i\beta} u_\beta$$

$$\frac{dp^0}{d\tau} = \frac{q}{c} F^{0\beta} u_\beta$$

$$\Leftrightarrow \frac{dE}{dt} = q \overset{\mathbf{r}}{E} \cdot \overset{\mathbf{r}}{v}$$

# Relativistic electrodynamics

$$\partial_{\alpha} F^{\alpha\beta} = \frac{4\pi}{c} J^{\beta}$$

$$\partial^{\alpha} F^{\beta\gamma} + \partial^{\beta} F^{\gamma\alpha} + \partial^{\gamma} F^{\alpha\beta} = 0$$

$$\frac{dp^{\alpha}}{d\tau} = \frac{q}{c} F^{\alpha\beta} u_{\beta}$$

Exactly the same as Maxwell's electrodynamics!

# Assignment

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1. Describe the transformation of force according to special theory of relativity.
2. Express Maxwell's equations in covariant form.
3. Describe the meaning of 4-vector notation.
4. Express Electromagnetic field in covariant form.



Thank You